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RESEARCH PAPER

Determination of the Astrophysical S-factor and Thermonuclear Reaction Rates of the (α,n) Medium Elements Reactions

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ABSTRACT:

Cross-sections of the (α ,n) medium elements reactions as a function of energies of alpha (α)-particle such as ⁴⁵Sc(α ,n)⁴⁸V, ⁴⁸Ti(α ,n)⁵¹Cr , ⁵¹V(α ,n)⁵⁴Mn , ⁵⁰Cr(α ,n)⁵³Fe , ⁵⁵Mn(α ,n)⁵⁸Co , ⁵⁴Fe(α ,n)⁵⁷Ni , ⁵⁹Co(α ,n)⁶²Cu , ⁶²Ni(α ,n)⁶⁵Zn , ⁶³Cu(α ,n)⁶⁶Ga , and ⁶⁶Zn(α ,n)⁶⁹Ge have been interpolated from threshold to 10 MeV in step of 0.05 MeV by using the Program of MATLAB. Weighted averages of the Cross-sections in (mb) have been utilized to calculate the astrophysical S-factor and thermonuclear reaction rates as a function of the energy of the center of mass, E_{c.m.} and T₉ Which is the temperature in units of 10⁹K (T₉ = 10⁻⁹T)

respectively. Polynomial relationships have been utilized to fit the computed astrophysical S-factor and thermonuclear reaction rates at various T₉ from best fitting equations with the minimum Chi-Square. Empirical formulae of set of reactions ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{54}Mn$, ${}^{55}Mn(\alpha,n){}^{58}Co$, ${}^{59}Co(\alpha,n){}^{62}Cu$, and ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{56}Mn$, ${}^{66}Zn(\alpha,n){}^{69}Ge$ have been utilized to compute astrophysical S-factor as a function of E_{c.m.} and Z and thermonuclear reaction rates as a function of T₉ and the target nucleus atomic number Z. The results have been compared with the embraced astrophysical S-factor and thermonuclear reaction rates that have been calculated from the fitting equations which have a good agreement.

KEY WORDS: Cross-sections; astrophysical S-factor; thermonuclear reaction rates; Gamow factor; Gamow energy; Sommerfeld parameter.

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INTRODUCTION :

The astrophysical S-factor, S(E), has covered a large area which used in the field to remove the energy dependence of the Coulomb barrier penetration from the cross-section, $\sigma(E)$ (Jose, 2016). As stellar energies are much lower than the Coulomb barrier, the cross sections hardly depend on energy (Descouvemont, 2011).

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Thermonuclear reactions play an important role in supplying the major source of energy in stars in particular during hydrogen burning. This burning process in the stellar interiors consists of the proton-proton (pp) chain and the carbonnitrogen-oxygen (CNO) cycle (Abdul Aziz, 2008). The quantity of interest in computing thermonuclear reaction rates for astrophysical aims is $N_A < \sigma v >$, which is the product of Avogadro's number with the average value of the cross section times velocity, averaged over a Maxwell-Boltzmann distribution of temperature (Roughton et al., 1983). Total Cross-sections of the (α,n) medium element reactions, that is a function of center of mass energy, have been calculated by a few authors, which are reminded by various references such as ${}^{45}Sc(\alpha,n){}^{48}V(Vlieks,$ Morgan and Blatt, 1974; Hansper et al., 1989; Haider, 2012), ${}^{48}\text{Ti}(\alpha,n){}^{51}\text{Cr}(\text{Chang et al., 1973};$ Vonach, Haight and Winkler, 1983; Levkovski, 1991; Morton et al., 1992; Baglin, Coral et al., 2004), ${}^{51}V(\alpha,n){}^{54}Mn$ (Levkovski, 1991; Hansper et al., 1993; Sonzogni et al., 1993; Peng, He and Long, 1999; Noori, 2008; Haider, 2012), 50 Cr(α ,n) 53 Fe(Vlieks, Morgan and Blatt, 1974; Morton *et al.*, 1994; Haider, 2012) ⁵⁵Mn(α,n)⁵⁸Co(Rizvi *et al.*, 1989; Levkovski, 1991; Tims et al., 1993; Haider, 2012) 54 Fe(α ,n)⁵⁷Ni (Houck and Miller, 1961; Vlieks, Morgan and Blatt, 1974; Tims et al., 1991; Haider, 2012) , ${}^{59}Co(\alpha,n){}^{62}Cu(Stelson and$ McGowan, 1964; D`auria et al., 1968; Zhukova et al., 1972; Tims et al., 1988; Noori, 2008) 62 Ni(α ,n) 65 Zn (Stelson and McGowan, 1964; Levkovski, 1991: Haider. 2012) 63 Cu(α ,n) 66 Ga(Stelson and McGowan, 1964; Zhukova et al., 1970; Haider, 2012), and 66 Zn(α ,n) 69 Ge(Stelson and McGowan, 1964: Levkovski, 1991) respectively. The goal of this work is to determine the empirical formulae to compute the astrophysical S-factor, S(E), and thermonuclear reaction rates, $N_A < \sigma v >$, utilizing the altered cross- sections of the reactions of the medium elements. The outcomes were compared with those published in the previous work.

2. Theory

Atomic masses of each medium element and isotopes related to this present work have been taken from the nuclear wallet cards published by the National Nuclear Data Center (NNDC) (Tuli, 2011). The *Q*-Value of the reaction $X(\alpha, n)Y$, is defined as the difference between the initial and the final rest mass energies (Meyerhof, 1967):

$$Q = [M_{\alpha} + M_X - (M_Y + M_n]c^2$$
(1)

Where $(M_{\alpha}, M_X, M_Y, and M_n)$ are the atomic masses of the incident, target particles, product nucleus and neutron (outgoing particle), respectively and $(c^2 = 931.494013 \text{ MeV/u}; \text{ where u=atomic mass unit (amu)} = 1.66 \times 10^{-27} \text{ kg})$. This equation is called the Q-value equation. If Q is + ive, the reaction called exoergic; if Q is - ive, it is endoergic.

The amount of energy needed for an endoergic reaction is called the *threshold energy* and can be calculated easily (Kaplan, 1962).

$$E_{th} = -Q\left(1 + \frac{M_{\alpha}}{M_X}\right) \tag{2}$$

Fusion requires two (or more) interacting particles to approach closely enough, within the short range of the (attractive) strong nuclear force, $\leq 10^{-15}$ m, to construct a new nucleus with A = A₁+A₂. The so-called height V_C of the barrier is its maximum value, which occurs at the nuclear radius, and is (Evans, 1955).

$$V_C = \frac{Z_1 Z_2 e^2}{R} \tag{3}$$

Where Z_1 and Z_2 are the charges of the projectile and target nuclei, and R and $(R = R_1 + R_2)$ is their separation, e is the charge of electron ($e^2 =$ 1.44 *MeV fm*), and the radius of the nucleus is given by $R = 1.3 \times 10^{-13} A^{1/3}$ cm, where A is the mass number (atomic weight) (Shaviv, 2012). Then Eq. (3) leads to

$$V_C = E_C = \frac{1.44}{1.3} \left(\frac{Z_1 Z_2}{A_1^{1/3} + A_2^{1/3}} \right) \tag{4}$$

Where E_c is the coulomb barrier (Coulomb energy) in MeV, $A_1^{1/3}$ and $A_2^{1/3}$ are the mass numbers of the charges of bombarding and targeting nuclei respectively.

The astrophysical S-factor, S(E), in the unit (*MeV-b*) is related to the cross-section by (Li, J. *et al.*, 2012):

$$S(E) = E\sigma(E) \exp(2\pi\eta)$$
(5)

Where *E* is the center-of-mass energy ($E_{c.m.}$) in *MeV*, $\sigma(E)$ is the cross-section of the reaction in (mb), $2\pi\eta$ is the Gamow factor, and η is Sommerfeld parameter (Angulo *et al.*, 1999):

$$\eta = \frac{Z_1 Z_2 e^2}{\hbar v} = 0.1575 Z_1 Z_2 \sqrt{\frac{\mu(u)}{E(MeV)}}$$
(6)

, \hbar is Planck's constant over $2\pi (1.0546 \times 10^{-27} \text{ ergs})$, ν is the relative velocity, μ is the reduced mass. The Gamow factor G (E) or $2\pi\eta$ can be written as in (Jose, 2016):

$$2\pi\eta = 0.98951 Z_1 Z_2 \sqrt{\frac{\mu(u)}{E(MeV)}}$$
(7)

The reduced mass μ in u (amu) is determined by the relationship (Clayton, 1968):

$$\mu = \frac{m_1 m_2}{m_1 + m_2} \tag{8}$$

Where m_1 and m_2 represent the masses of the bombarding and target nucleus in units of (amu), respectively. The energy of the center of mass of pair of particles $E_{c.m.}$ is related to the laboratory energy, $E_{Lab.}$ of the projectile particle by the equation (Meyerhof, 1967):

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$$E_{c.m.} = \frac{m_2}{m_1 + m_2} E_{lab.}$$
(9)

The Gamow energy E_G , in *MeV* (Brown, 2015):

$$E_G = 2\pi^2 \mu u C^2 \alpha^2 (Z_1 Z_2)^2 = 0.979 \mu (Z_1 Z_2)^2$$
(10)

Where $\alpha = \frac{1}{137} = \frac{e^2}{\hbar c}$ is the fine-structure constant. The thermonuclear reaction rates, $N_A \langle \sigma v \rangle$ in unit $(cm^3mol^{-1}s^{-1})$ (Angulo *et al.*, 1999):

$$N_A \langle \sigma v \rangle = \left(\frac{8}{\mu \pi}\right)^{1/2} \frac{1}{(k_B T)^{3/2}} N_A \int_0^\infty E \sigma(E) \exp(-E/k_B T) dE$$
(11)

Where N_A is the Avogadro's number ($6.022 \times 10^{23} mol^{-1}$), k_B is the Boltzmann's constants ($1.38 \times 10^{-16} erg/K$), and *T* is the temperature respectively. Eq. (11) leads to (Angulo *et al.*, 1999):

$$\begin{split} N_A \langle \sigma v \rangle &= 3.7313 \times 10^7 \mu^{-1/2} T_9^{-3/2} \int_0^\infty E \sigma(E) \exp(-11.605 E \\ &/T_9) \ dE \ (12) \end{split}$$

Where T_9 is the temperature in units of $10^9 K$ ($T_9 = 10^{-9}T$)

The weighted averages of the Cross-sections of medium elements $\sigma_0(mb)$ and the uncertainty (errors) $\Delta \sigma_0(mb)$ are expressed by the following Eqs. (Bevington and Robinson, 2003):

$$\sigma_0(mb) = \frac{\sum_i (\sigma_i / \delta_i^2)}{\sum_i (1 / \delta_i^2)}$$
(13)

Where σ_i and $\delta_i (\Delta \sigma_i)$ are the cross-section and the uncertainties of i^{th} reference, relating to each value of σ_i ,

$$\Delta\sigma_0(mb) = \pm \frac{1}{\sqrt{\sum_i (1/\delta_i^2)}} \tag{14}$$

The considered formalism type is the polynomial fit expression of the shape:

$$Y = C_0 + C_1 X + C_2 X^2 + C_3 X^3 + \dots + C_N X^N = \sum_{i=0}^M C_i X^i \quad (15)$$

This polynomial is obtained by the Excel computer program (Format Trendline). Where $(C_0, C_1, C_2, C_3, ...)$ are free parameters (coefficients of polynomial), and (i = 0, 1, 2, 3, ..., M), and

$$C_i = \sum_{j=0}^N C_{ij} K^j \tag{16}$$

Are considered in this work, then by combining the Eqs. (15) & (16), the following relation has been acquired:

$$Y = \sum_{i=0}^{M} (\sum_{j=0}^{N} C_{ij} K^{j}) X^{i}$$
(17)

Where Y=ln[S(E)] or $ln[N_A < \sigma v >]$, (i=0,1,2,...M), (j=0,1,2,...N), (C₀₀,C₀₁,C₀₂,...) are coefficients of polynomials, K is the energy of the center of mass

or T₉ according to the S(E) or N_A< σ v>, and X is atomic number Z. The Excel computer program has been utilized to acquire the best fit relationship corresponding to various energies ranges near threshold up to 10 *MeV* in the center of mass system or T₉ ranges from (1 to 10) 10⁹K. The data of these extents were avoided in each step, till a possible value of the determination coefficient $R^2 \approx 1$ was come to. The best fit adopted data was acquired with increasing order to supply the minimum value of Chi-Square (χ^2) by using the Eq. (Belgaid *et al.*, 2005):

$$\chi^{2} = \frac{1}{(N-M)} \sum_{i}^{N} \left(\frac{Y_{exp}^{i} - Y_{cal}^{i}}{\Delta Y_{exp}^{i}} \right)^{2}$$
(18)

Where *N* is the data points' number, *M* is the fitting coefficients number, Y_{exp}^i and ΔY_{exp}^i are the experimental (adopted value) of ln[S(E)] or $ln[N_A(\sigma v)]$ and its error (uncertainty) respectively, Y_{cal}^i is the calculated ln[S(E)] or $ln[N_A(\sigma v)]$.

3. Data Reduction and Analysis

The Atomic masses have been taken into consideration to determine the Q-Value, threshold energy, Coulomb barrier, reduced mass, and the ratio between $(E_{c.m}/E_{lab.})$ of (α, n) medium elements reactions using the Eqs. (1, 2, 4, 8, and 9); the results have been shown in the table (1). Eqs. (6,7,10, and 5) taken into consideration to

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determine the Sommerfeld parameter(η), Gamow factor G(E), Gamow energy (E_G), and the S-factor of astrophysical, S(E) of the (α ,n) medium element reactions. The results are shown in table (2). The cross-sections of (α ,n) reactions of medium elements in present work such as (⁴⁵Sc, ⁴⁸Ti, ⁵¹V, ⁵⁰Cr, ⁵⁵Mn, ⁵⁴Fe, ⁵⁹Co, ⁶²Ni, ⁶³Cu, and ⁶⁶Zn), which are available in the literature review has been taken and plotted again, and using the MATLAB software to interpolate to acquire the cross-sections in fine step of 0.05 MeV. The weighted average of the altered Cross-sections of ith references for the medium elements which cross-section (σ_0) and uncertainty ($\Delta \sigma_0$) have been computed by using Eqs. (13) and (14)

respectively. The acquired results have been utilized to calculate the astrophysical S-factor and thermonuclear reaction rates of (α ,n) reactions as a function of the center of mass energies $E_{c.m.}$ by using eq. (5) and (12). The acquired equations to compute the S-factor of the reminded reactions are shown in Table 2. The final formula for each astrophysical S-factor,

The final formula for each astrophysical S-factor, S(E) and thermonuclear reaction rates $N_A \langle \sigma v \rangle$ is shown in Eq. (17) where $Y = \ln[S(E)]$ or $Y = \ln[N_A \langle \sigma v \rangle]$.

Table 1. Q-Value, threshold energy ($E_{threshold}$), Coulomb barrier E_c , reduced mass (μ), and the ratio between

| (α,n) | O-value | E threshold | (MeV) | Coulomb Barrier | Reduced | |
|--|--|-------------|-----------------------|-----------------|-----------|-------------|
| Medium Element | (MeV) | | G 1 G 1 | E (MeV) | Mass (µ) | Ec.m./Elab. |
| Reaction | $(\mathbf{W}\mathbf{I}\mathbf{C}\mathbf{V})$ | Lab. System | C.M. System | L_c (ivic v) | (amu) | |
| ${}^{45}Sc(\alpha,n){}^{48}V$ | -2.241E+00 | 2.440E+00 | 2.241E+00 | 9.044E+00 | 3.675E+00 | 9.182E-01 |
| ${}^{48}\text{Ti}(\alpha,n){}^{51}\text{Cr}$ | -2.687E+00 | 2.911E+00 | 2.687E+00 | 9.334E+00 | 3.694E+00 | 9.230E-01 |
| $^{51}V(\alpha,n)^{54}Mn$ | -2.294E+00 | 2.474E+00 | 2.294E+00 | 9.622E+00 | 3.711E+00 | 9.272E-01 |
| ${}^{50}Cr(\alpha,n){}^{53}Fe$ | -4.961E+00 | 5.359E+00 | 4.961E+00 | 1.009E+01 | 3.706E+00 | 9.258E-01 |
| ⁵⁵ Mn(α,n) ⁵⁸ Co | -3.512E+00 | 3.767E+00 | 3.512E+00 | 1.027E+01 | 3.731E+00 | 9.321E-01 |
| 54 Fe(α ,n) 57 Ni | -5.817E+00 | 6.249E+00 | 5.817E+00 | 1.073E+01 | 3.726E+00 | 9.309E-01 |
| 59 Co(α ,n) 62 Cu | -5.089E+00 | 5.434E+00 | 5.089E+00 | 1.091E+01 | 3.748E+00 | 9.364E-01 |
| 62 Ni(α ,n) 65 Zn | -6.480E+00 | 6.899E+00 | 6.480E+00 | 1.119E+01 | 3.760E+00 | 9.393E-01 |
| 63 Cu(α ,n) 66 Ga | -7.502E+00 | 7.979E+00 | 7.502E+00 | 1.154E+01 | 3.763E+00 | 9.402E-01 |
| ${}^{66}Zn(\alpha,n){}^{69}Ge$ | -7.445E+00 | 7.897E+00 | 7.445E+00 | 1.181E+01 | 3.774E+00 | 9.428E-01 |

 $(E_{c.m./Elab.})$ of (α, n) medium elements reactions.

Table 2. The Sommerfeld parameter(η), Gamow factor G(E), Gamow energy (E_G), and the astrophysical Sfactor, S(E) of the (α ,n) medium elements reactions

| (α,n) Medium Element Reaction | Sommerfeld Parameter η | Gamow factor G(E) | Gamow Energy Eg(MeV) | Astrophysical S-factor S(E) |
|--|------------------------------|-----------------------------|-------------------------|--|
| $^{45}Sc(\alpha,n)^{48}V$ | $1.268E+01/\sqrt{E_{c.m.}}$ | $7.967E+01/\sqrt{E_{c.m.}}$ | 6.348E+03 | $E_{c.m.}\sigma(E)Exp(7.967E+01/\sqrt{E_{c.m.}})$ |
| $^{48}\text{Ti}(\alpha,n)^{51}\text{Cr}$ | $1.331E+01/\sqrt{E_{c.m.}}$ | $8.368E+01/\sqrt{E_{c.m.}}$ | 7.003E+03 | $E_{c \cdot m} \sigma(E) Exp(8.368E+01/\sqrt{E_{c.m.}})$ |
| ${}^{51}V(\alpha,n){}^{54}Mn$ | $1.395E+01/\sqrt{E_{c.m.}}$ | $8.769E+01/\sqrt{E_{c.m.}}$ | 7.689E+03 | $E_{c \cdot m} \sigma(E) Exp(8.769E+01/\sqrt{E_{c.m.}})$ |
| ${}^{50}Cr(\alpha,n){}^{53}Fe$ | $1.455E+01/\sqrt{E_{c.m.}}$ | $9.143E+01/\sqrt{E_{c.m.}}$ | 8.360E+03 | $E_{c \cdot m} \sigma(E) Exp(9.143E+01/\sqrt{E_{c.m.}})$ |
| 55 Mn(α ,n) 58 Co | $1.520E+01/\sqrt{E_{c.m.}}$ | $9.556E+01/\sqrt{E_{c.m.}}$ | 9.132E+03 | $E_{c \cdot m} \sigma(E) Exp(9.556E+01/\sqrt{E_{c.m.}})$ |
| ${}^{54}\text{Fe}(\alpha,n){}^{57}\text{Ni}$ | $1.580E+01/\sqrt{E_{c.m.}}$ | $9.932E+01/\sqrt{E_{c.m.}}$ | 9.865E+03 | $E_{c.m.}\sigma(E)Exp(9.932E+01/\sqrt{E_{c.m.}})$ |
| 59 Co(α ,n) 62 Cu | $1.646E+01/\sqrt{E_{c.m.}}$ | $1.034E+02/\sqrt{E_{c.m.}}$ | 1.070E+04 | $E_{c \cdot m} \sigma(E) Exp(1.034E+02/\sqrt{E_{c.m.}})$ |
| 62 Ni(α ,n) 65 Zn | $1.709E+01/\sqrt{E_{c.m.}}$ | $1.074E+02/\sqrt{E_{c.m.}}$ | 1.154E+04 | $E_{c \cdot m} \sigma(E) Exp(1.074E+02/\sqrt{E_{c.m.}})$ |
| $^{63}Cu(\alpha,n)^{66}Ga$ | $1.771E+01/\sqrt{E_{c.m.}}$ | $1.113E+02/\sqrt{E_{c.m.}}$ | 1.240E+04 | $E_{c.m.}\sigma(E)Exp(1.113E+02/\sqrt{E_{c.m.}})$ |
| 66 Zn(α ,n) 69 Ge | $1.835E+01/\sqrt{E_{c.m.}}$ | $1.153E+02/\sqrt{E_{c.m.}}$ | 1.330E+04 | $E_{c \cdot m} \sigma(E) Exp(1.153E + 02/\sqrt{E_{c.m.}})$ |

4. Results and Discussion

In general, we can write Eq. (15), and instead of *X* insert center of mass energies $E_{c.m.}$. Then the Eq. (15) becomes

$$Y = C_0 + C_1 K + C_2 K^2 + C_3 K^3 + \dots + C_N K^N$$
$$= \sum_{i=0}^{M} C_i K^i$$
(19)

Where (C₀, C₁, C₃...) are free parameters, K are parameters that represent the C.M energy or T₉, (i=0, 1, 2, 3... M), and Y=ln[S-factor (MeV-b)] or Y=ln[N_A< σ v> (cm³mol⁻¹s⁻¹)].

4.1. Astrophysical S-factor Empirical Formulae The adopted astrophysical S-factor has been used to acquire the fitting parameters by using the expressions of the polynomial (18), (20) and (19) as shown in the steps:

1. The polynomial relations which are utilized in eq. (19) to fit the computed astrophysical S-factor, S(E) in the natural logarithm of the calculated elements to compute the adopted (taken on) natural logarithm of astrophysical S-factor from the best fitting with a minimum (χ^2) using Eq. (20). The acquired best fitting relations of the reminded reactions were presented in Eqs. (20, 21, 22, 23, 24, 25, 26, 27, 28, and 29) for the reactions ⁴⁵Sc(α ,n)⁴⁸V, ⁴⁸Ti(α ,n)⁵¹Cr, ⁵¹V(α ,n)⁵⁴Mn, ⁵⁰Cr(α ,n)⁶³Fe, ⁵⁵Mn(α ,n)⁵⁸Co, ⁵⁴Fe(α ,n)⁵⁷Ni, ⁵⁹Co(α ,n)⁶²Cu, ⁶²Ni(α ,n)⁶⁵Zn, ⁶³Cu(α ,n)⁶⁶Ga, and ⁶⁶Zn(α ,n)⁶⁹Ge respectively.

 ${}^{45}Sc(\alpha,n)^{48}V \qquad x^2 = 0.0247 \\ ln[S - factor(MeV - b)] = 0.0062E^3 - 0.2075E^2 + \\ 1.0183E + 30.7 \tag{20}$

⁴⁸Ti(α ,n)⁵¹Cr $x^2 = 0.086$ ln[S - factor (MeV - b)] = -0.0867E⁴ + 2.5599E³ - 28.125E² + 135.15E - 206.65 (21)

 ${}^{51}V(\alpha,n)^{54}Mn \qquad x^2 = 0.041$ $ln[S - factor (MeV - b)] = -0.0261E^3 + 0.4478E^2 - 3.1956E + 41.906$ (22)

 ${}^{55}Mn(\alpha,n)^{58}Co \qquad x^2 = 0.0015$ $ln[S - factor (MeV - b)] = 0.0153E^3 - 0.4515E^2 + 3.3053E + 28.436$ (24)

 62 Ni(α ,n) 65 Zn $x^2 = 0.0056$

 $ln[S - factor (MeV - b)] = 0.1946E3E^{3} - 4.9013E^{2} + 40.02E - 69.093$ (27)

 ${}^{63}Cu(\alpha,n)^{66}Ga \qquad x^2 = 0.027$ $ln[S - factor (MeV - b)] = 2.2587E^3 - 57.223E^2 + 480.3E - 1299.5 \qquad (28)$

2. At fixed values of energy in center-of-mass, the change of the S-factor in natural logarithm with the Z has been fitted to the polynomial relation utilizing Eq. (19). The acquired results were used to determine the free parameters (coefficients of polynomial) (C_i).

3. The free parameters C_i , were plotted against each value of the center of mass energies and fitted to sufficient the polynomial relation were shown in Eq. (16).

4. The last formula of a set of reactions has been calculated by utilizing the combination of the two polynomials to show the systematic manner of the reactions which are shown in Eq. (17). The Y Variable is the astrophysical S-factor.

4.1.1 The Empirical Formulae Relating the Astrophysical S-factor to Center of

Mass Energy and the Atomic Number Z of the Target Nucleus

The empirical formulae related to the astrophysical S-factor (MeV-b) with both of center of mass energy E_{c.m.}, and the atomic number Z were performed as the steps below: 1- At fixed values of the center of mass energies from 5.5 to 10 MeV in steps of 0.25 MeV for the ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{54}Mn$ $^{55}Mn(\alpha,n)^{58}Co$, and $^{59}Co(\alpha,n)^{62}Cu$ reactions, the astrophysical S- factor in natural logarithm will vary with the atomic number(Z), as shown in Fig. (1). The data was fitted into the accompanying polynomial expression:

$$Y = \sum_{i=0}^{2} C_i X^i \tag{30}$$

Where $Y = \ln[S(E)]$, and X=Z, with free parameters C_i (C₀, C₁, and C₂).

2- The S-factor, S(E), which is was adopted, has been utilized as a function of atomic number Z of target nucleus at the fixed center of mass energies using the computer program Excel to acquire the fitting relations and then it was utilized to compute the fitting parameters. The acquired results are shown in Table 3.

3- The obtained free parameters C_i (C_0 , C_1 , and C_2), presented in Table (3) are plotted against with the fixed values of center of mass energies from 5.5 to 10 MeV in step of 0.25 MeV as shown in Fig.(2), and then the acquired coefficients of polynomials C_i have been fitted to the polynomial expression below:

$$C_i = \sum_{j=0}^{L} C_{ij} E^j \tag{31}$$

The combination of the two polynomials Eq. (30) and Eq. (31) takes the form of the formula below of energy ranged from 5.5 to 10 MeV in the step of 0.25 MeV:

$$Y = \sum_{i=0}^{2} (\sum_{j=0}^{2} C_{ij} E^{j}) X^{i} \qquad (32)$$

Where Y=ln[S(E)], X=atomic number Z

$$Y = \sum_{i=0}^{2} (C_{i0}E^{0} + C_{i1}E^{1} + C_{i2}E^{2})X^{i}$$

$$Y = C_{00}E^{0}X^{0} + C_{01}E^{1}X^{0} + C_{02}E^{2}X^{0} + C_{10}E^{0}X^{1} + C_{11}E^{1}X^{1} + C_{12}E^{2}X^{1} + C_{20}E^{0}X^{2} + C_{21}E^{1}X^{2} + C_{22}E^{2}X^{2}$$
(33)

Where $(C_{00}, C_{01}, C_{02}, C_{10}, C_{11}, \dots, C_{22})$ are free parameters and their values are shown in the matrix below:

$$\begin{bmatrix} C_{00} & C_{01} & C_{02} \\ C_{10} & C_{11} & C_{12} \\ C_{20} & C_{21} & C_{22} \end{bmatrix} = \begin{bmatrix} -2.0079 & -14.163 & 1.4295 \\ 2.5083 & 1.1719 & -0.1305 \\ -0.0466 & -0.0217 & 0.0027 \end{bmatrix}, \begin{bmatrix} R^2 = 0.6757 \\ R^2 = 0.7426 \\ R^2 = 0.7585 \end{bmatrix}$$

The acquired formula of a set of reactions such as ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{54}Mn$, ${}^{55}Mn(\alpha,n){}^{58}Co$, and ${}^{59}Co(\alpha,n){}^{62}Cu$ has been used to calculate the astrophysical S-factor S(E) for each of the above reactions and compared with the adopted astrophysical S-factor calculated from the fitting expressions and shown to be in a good agreement and the comparison of the two results are shown in Table (4).

Table 3. Free parameters C_i (C_0 , C_1 , and C_2) as a function of the energy of the center of mass .

| Ec.m. (MeV) | C0 | C1 | C2 |
|----------------|---------|---------|---------|
| 5.5 | -49.997 | 6.1814 | -0.1106 |
| 5.75 | -41.303 | 5.3931 | -0.093 |
| 6 | -34.449 | 4.7574 | -0.0787 |
| 6.25 | -29.331 | 4.2685 | -0.0675 |
| 6.5 | -25.812 | 3.9172 | -0.0593 |
| 6.75 | -23.719 | 3.6906 | -0.0538 |
| 7 | -22.843 | 3.5724 | -0.0508 |
| 7.25 | -22.944 | 3.5423 | -0.0497 |
| 7.5 | -23.743 | 3.5767 | -0.0501 |
| 7.75 | -24.928 | 3.6484 | -0.0515 |
| 8 | -26.154 | 3.7264 | -0.053 |
| 8.25 | -27.038 | 3.7763 | -0.054 |
| 8.5 | -27.164 | 3.7599 | -0.0535 |
| 8.75 | -26.081 | 3.6355 | -0.0508 |
| 9 | -23.304 | 3.3579 | -0.0446 |
| 9.25 | -18.311 | 2.8781 | -0.0339 |
| 9.5 | -10.548 | 2.1436 | -0.0176 |
| 9.75 | 0.5769 | 1.0983 | 0.0057 |
| 10 | 15.687 | -0.3175 | 0.0374 |

Table 4. Comparison between polynomial fitting expression (Best Fitting) of the adopted astrophysical S-Factor of (α,n) medium element reactions with those computed from Eq. (33).

| Ec.m | $^{45}\mathrm{Sc}(\alpha,n)^{48}\mathrm{V}$ | ${}^{48}\mathrm{Ti}(\alpha,n){}^{51}\mathrm{C}$ | 51 V(α ,n) 54 Mn | $^{55}Mn(\alpha,n)^{58}Co$ | 59 Co(α ,n) 62 Cu | |
|------|---|---|------------------------------------|----------------------------|-------------------------------------|--|
| | | | | | | |

| 70 | |
|----|--|
| 10 | |

| | ln[S- | 1 | ln[S- | ln[S- | ln[S- | ln[S- | ln[S- | ln[S- | ln[S- | |
|-------|-------------|-------------------|-----------|---------|----------|---------|-------------------|---------|---------------------|-----------|
| (Me | factor(Me | In[S- factor(M | factor(Me | factor(| factor(M | factor(| factor(M | factor(| factor(MeV | ln[S- |
| V) | V-b)] | Tactor(M | V-b)] | MeV- | eV-b)] | MeV- | eV-b)] | MeV- | -b)] | factor(Me |
| | (Best | ev-b)] | (Best | b)] | (Best | b)] | (Best | b)] | (Best | V-b)] |
| | Fitting) | (Formula | Fitting) | (Form | Fitting) | (Formul | Fitting) | (Formul | Fitting) | (Formula) |
| | 4.04% |) | 4.215% | ula) | 3.801% | a) | 3.028% | a) | 1.974% | |
| | 31.055+1. | | 32.461+1. | | 33,534+1 | | 35,503+1 | | 36.214+0.7 | |
| 5.5 | 255 | 31.301 | 368 | 32.684 | 275 | 33.897 | 075 | 35.819 | 15 | 37.067 |
| | 20 972+1 | | 22 466+1 | | 22 275+1 | | 25 422+1 | | 26 472+0 7 | |
| 5.75 | 50.875±1. | 31.182 | 32.40011. | 32.584 | 33.37311 | 33.821 | 55.422 <u>1</u> 1 | 35.803 | 30.473±0.7 | 37.128 |
| | 247 | | 308 | | .269 | | .073 | | 20 | |
| 6 | 30.679±1. | 31.048 | 32.325±1. | 32.467 | 33.216±1 | 33.727 | 35.319±1 | 35.769 | 36.645±0.7 | 37.174 |
| - | 239 | | 363 | | .263 | | .069 | | 23 | - |
| 6.25 | 30.473±1. | 30 800 | 32.087±1. | 32 331 | 33.054±1 | 33 615 | 35.193±1 | 35 716 | 36.737±0.7 | 37 203 |
| 0.25 | 231 | 50.055 | 352 | 52.554 | .256 | 55.015 | .066 | 55.710 | 25 | 57.205 |
| 65 | 30.255±1. | 20 725 | 31.791±1. | 22 102 | 32.886±1 | 22.404 | 35.046±1 | | 36.758±0.7 | 27 217 |
| 0.5 | 222 | 30.735 | 340 | 32.183 | .250 | 33.484 | .061 | 35.645 | 26 | 37.217 |
| | 30.026±1. | | 31.472±1. | | 32.712±1 | | 34.881±1 | | 36.715±0.7 | |
| 6.75 | 213 | 30.556 | 327 | 32.016 | 243 | 33.336 | 056 | 35.556 | 25 | 37.216 |
| | 213 | | 21 15/1+1 | | 22 527+1 | | 24 609±1 | | 26 619+0 7 | |
| 7 | 29.767±1. | 30.362 | 51.154±1. | 31.832 | 52.527±1 | 33.170 | 54.096±1 | 35.449 | 50.010±0.7 | 37.198 |
| | 203 | | 313 | | .230 | | .051 | | 23 | |
| 7.25 | 29.539±1. | 30 152 | 30.853±1. | 31 631 | 32.329±1 | 32 986 | 34.498±1 | 35 324 | 36.474±0.7 | 37 165 |
| | 193 | 00.202 | 300 | 01.001 | .229 | 01000 | .045 | | 20 | 0/1200 |
| 75 | 29.281±1. | 20 028 | 30.577±1. | 21 /12 | 32.117±1 | 22 781 | 34.284±1 | 25 120 | 36.292±0.7 | 27 117 |
| 7.5 | 183 | 29.920 | 289 | 51.415 | .221 | 52.764 | .038 | 55.100 | 16 | 57.117 |
| | 29.015±1. | 20,000 | 30.328±1. | 24.470 | 31.887±1 | 22 564 | 34.056±1 | 25.040 | 36.079±0.7 | 27.052 |
| 1.15 | 172 | 29.689 | 278 | 31.179 | .212 | 32.564 | .031 | 35.019 | 12 | 37.052 |
| | 28.741+1. | | 30.096+1. | | 31.637+1 | | 33.816+1 | | 35.844+0.7 | |
| 8 | 161 | 29.434 | 269 | 30.928 | 203 | 32.326 | 024 | 34.839 | 08 | 36.972 |
| | 20 / 50 + 1 | | 205 | | 21 265+1 | | 22 566+1 | | 25 505+0 7 | |
| 8.25 | 20.439±1. | 29.165 | 29.80311. | 30.659 | 51.505±1 | 32.070 | 55.500±1 | 34.641 | 55.595 <u>+</u> 0.7 | 36.876 |
| | 150 | | 259 | | .192 | | .010 | | 03 | |
| 8.5 | 28.171±1. | 28.880 | 29.613±1. | 30.374 | 31.068±1 | 31.796 | 33.306±1 | 34.424 | 35.340±0.6 | 36.764 |
| | 138 | | 248 | | .181 | | .009 | - | 98 | |
| 8 75 | 27.877±1. | 28 581 | 29.306±1. | 30 072 | 30.744±1 | 31 50/ | 33.039±1 | 3/1 100 | 35.088±0.6 | 36 637 |
| 0.75 | 126 | 28.581 | 235 | 30.072 | .169 | 51.504 | .000 | 54.150 | 93 | 30.037 |
| 0 | 27.577±1. | 20.200 | 28.903±1. | 20 75 4 | 30.391±1 | 24.405 | 32.766±0 | 22.027 | 34.846±0.6 | 26.404 |
| 9 | 114 | 28.266 | 218 | 29.754 | .155 | 31.195 | .992 | 33.937 | 88 | 36.494 |
| | 27 272+1 | | 28 357+1 | | 30 005+1 | | 32 488+0 | | 34 622+0 6 | |
| 9.25 | 102 | 27.937 | 195 | 29.418 | 140 | 30.867 | 984 | 33.666 | 83 | 36.335 |
| | 26.062+1 | | 27 611 1 | | 20 59411 | | 22.20610 | | 24.42610.6 | |
| 9.5 | 20.903±1. | 27.592 | 27.011±1. | 29.066 | 29.584±1 | 30.521 | 32.206±0 | 33.377 | 34.420±0.0 | 36.161 |
| | 089 | | 104 | | .124 | | .975 | | 80 | |
| 9.75 | 26.649±1. | 27.232 | 26.600±1. | 28,696 | 29.127±1 | 30,157 | 31.923±0 | 33.070 | 34.264±0.6 | 35.971 |
| 1.1.0 | 077 | | 121 | | .107 | 50.207 | .967 | | 76 | 00.072 |
| 10 | 26.333±1. | 26.858 | 25.250±1. | 28 210 | 28.630±1 | 20 776 | 31.639±0 | 22 7/5 | 34.146±0.6 | 35 765 |
| 10 | 064 | 20.030 | 064 | 20.310 | .088 | 23.770 | .958 | 52.745 | 74 | 33.703 |



Fig. 1. The variation of the natural logarithm of the astrophysical S-factor S(E) with the atomic number (Z) for the ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{54}Mn$, ${}^{55}Mn(\alpha,n){}^{58}Co$, and ${}^{59}Co(\alpha,n){}^{62}Cu$ reactions at fixed values of center of mass energies.



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Fig. 2. C_i coefficients against the center of mass energy, for C_0 , C_1 , and C_2 respectively. The solid line represents the fitted curve through the data.

4.2. Thermonuclear Reaction Rates Empirical Formulae

The adopted thermonuclear reaction rates $N_A < \sigma v >$ have been utilized to acquire the fitting parameter by utilizing the polynomial expressions (16), (18) and (19) by the steps below:

1. Polynomial expressions were utilized in eq. (19) to fit the computed thermonuclear reaction rates natural logarithm $N_A < \sigma v >$ of the thoughtful medium elements to set the embraced natural logarithm of thermonuclear reaction rates $N_A < \sigma v >$ from the best fitting with a minimum (χ^2) utilizing Eq. (18). The acquired best fitting relationships of the remembered reactions are shown in Eqs. (34, 35, 36, 37, 38, 39, 40, 41, 42, and 43) for the reactions ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{54}Mn$, ${}^{50}Cr(\alpha,n){}^{53}Fe$, ${}^{55}Mn(\alpha,n){}^{58}Co$, ${}^{54}Fe(\alpha,n){}^{57}Ni$, ${}^{59}Co(\alpha,n){}^{62}Cu$, ${}^{62}Ni(\alpha,n){}^{65}Zn$, ${}^{63}Cu(\alpha,n){}^{66}Ga$, and $^{66}Zn(\alpha,n)^{69}Ge$ respectively. 45 Sc(α ,n) 48 V $x^2 = 1.064$ $\ln[N_A \langle \sigma v \rangle (cm^3 \, s^{-1} \, mol^{-1})] = 0.0018T^5 - 0.0636T^4 +$ $0.9094T^3 - 6.7134T^2 + 27.463T - 42.483$ (34) 48 Ti(α ,n) 51 Cr $x^2 = 1.216$ $5.103T^2 + 26.92T - 48.328$ (35) $^{51}V(\alpha,n)^{54}Mn$ $x^2 = 0.979$ $ln[N_{A}(\sigma v)(cm^{3} s^{-1} mol^{-1})] = -0.01T^{4} + 0.299T^{3} 3.4636T^2 + 19.663T - 37.011$ (36) ${}^{50}Cr(\alpha,n){}^{53}Fe$ $x^2 = 0.445$

 $ln[N_{A}(\sigma v)(cm^{3} s^{-1} mol^{-1})] = 0.0281T^{3} - 0.8253T^{2} + 8.9513T - 24.344$ (37) ⁵⁵Mn(a,n)⁵⁸Co x²=0.318 $ln[N_{A}(\sigma v)(cm^{3} s^{-1} mol^{-1})] = -0.0053T^{4} + 0.1798T^{3} - 2.4157T^{2} + 16.122T - 34.73$ (38) ⁵⁴Fe(a,n)⁵⁷Ni x²=0.57 $ln[N_{A}(\sigma v)(cm^{3} s^{-1} mol^{-1})] = -0.0093T^{4} + 0.2984T^{3} - 3.7055T^{2} + 22.31T - 47.654$ (39) ⁵⁹Co(a,n)⁶²Cu x²=0.849 $ln[N_{A}(\sigma v)(cm^{3} s^{-1} mol^{-1})] = 0.0013T^{5} - 0.0504T^{4} + 0.0504T^$

 $ln[N_{A}(\sigma v)(cm^{3} s^{-1} mol^{-1})] = 0.0013T^{5} - 0.0504T^{4} + 0.7943T^{3} - 6.5501T^{2} + 30.053T - 55.065$ (40)

2. At fixed values of T_9 , the variation of the natural logarithm of the thermonuclear reaction rates with the physical parameter atomic number Z has been fitted to the polynomial expression utilizing Eq. (19). The acquired outcomes are contemplated to set the coefficients of polynomials (C_i).

3. The coefficients of polynomials Ci_{τ} are plotted versus each value of T_9 and fitted to satisfactory

the polynomial expression were shown in Eq. (16).

4. The last formula of a set of reactions has been determined by utilizing the combination of the two polynomials to show the systematic manner of the reactions, which is shown in Eq. (17). The Y Variable is the thermonuclear reaction rates.

4.2.1. The Empirical Formulae Relating the Thermonuclear Reaction Rates to T₉ and the Atomic Number Z of the Target Nucleus

The empirical formulae relating to the thermonuclear reaction rates $N_A < \sigma v > (cm^3 s^{-1} mol^{-1})$ with both T₉ and Z were performed as the steps below:

1- At fixed values of the T₉ from 6 to 10 10⁹ K in steps of 0.25 10⁹ K for the ⁴⁵Sc(α ,n)⁴⁸V , ⁴⁸Ti(α ,n)⁵¹Cr , ⁵¹V(α ,n)⁵⁴Mn , ⁵⁵Mn(α ,n)⁵⁸Co , ⁶²Ni(α ,n)⁶⁵Zn , and ⁶⁶Zn(α ,n)⁶⁹Ge reactions, the natural logarithm of the thermonuclear reaction rates will vary with the atomic number Z this shown in Fig. (3). The data fitted to the polynomial expression as the same as Eq. (30), Where $Y = \ln[N_A \langle \sigma v \rangle]$, X=Z, with free parameters C_i (C₀, C₁, and C₂).

2- The adopted thermonuclear reaction rates have been used as a function of Z at fixed T_9 utilizing the computer program Excel to acquiring the fitting expressions and then used to calculate the fitting parameters. The obtained results are presented in Table (5).

3- The obtained free parameters C_i (C_0 , C_1 , and C_2), as presented in Table (5) are plotted versus with the fixed values of T_9 from 6 to 10 10^9 K in steps of 0.25 10^9 K as presented in Fig.(4),and then the acquired coefficients of polynomials C_i have been fitted to the polynomial expression:

$$C_i = \sum_{j=0}^{2} C_{ij} T_9{}^j \qquad (44)$$

The combination of the two polynomials Eq. (30) and Eq. (44) takes the shape of the following formula range T_9 from 6 to 10 10^9 K in steps of 0.25 10^9 K:

$$Y = \sum_{i=0}^{2} \left(\sum_{j=0}^{2} C_{ij} T_{9}^{j} \right) X^{i}$$
 (45)

Where $Y=\ln[N_A < \sigma v >]$, T₉ is the temperature in 10^9 K, and X=atomic numberZ

$$Y = \sum_{i=0}^{2} (C_{i0}T_{9}^{0} + C_{i1}T_{9}^{1} + C_{i2}T_{9}^{2})X^{i}$$

$$Y = C_{00}T_{9}^{0}X^{0} + C_{01}T_{9}^{1}X^{0} + C_{02}T_{9}^{2}X^{0} + C_{10}T_{9}^{0}X^{1} + C_{11}T_{9}^{1}X^{1} + C_{12}T_{9}^{2}X^{1} + C_{20}T_{9}^{0}X^{2} + C_{21}T_{9}^{1}X^{2} + C_{22}T_{9}^{2}X^{2}$$
(46)

Where $(C_{00}, C_{01}, C_{02}, C_{10}, C_{11}, \dots, C_{22})$ are free parameters and their values are shown in the matrix below:

$$\begin{bmatrix} C_{00} & C_{01} & C_{02} \\ C_{10} & C_{11} & C_{12} \\ C_{20} & C_{21} & C_{22} \end{bmatrix} =$$

$$\begin{bmatrix} 42.691 & -16.998 & 1.7602 \\ -3.0522 & 1.4998 & -0.145 \\ 0.0375 & -0.0265 & 0.0027 \end{bmatrix}, \begin{bmatrix} R^2 = 0.9717 \\ R^2 = 0.9645 \\ R^2 = 0.9643 \end{bmatrix}$$

The acquired formula of a set of reactions such as ${}^{45}Sc(\alpha,n){}^{48}V$, ${}^{48}Ti(\alpha,n){}^{51}Cr$, ${}^{51}V(\alpha,n){}^{54}Mn$, ${}^{55}Mn(\alpha,n){}^{58}Co$, ${}^{62}Ni(\alpha,n){}^{65}Zn$, and ${}^{66}Zn(\alpha,n){}^{69}Ge$ has been used to calculate the thermonuclear reaction rates $N_A\langle\sigma\nu\rangle$ for each of the above reactions and compared with the adopted thermonuclear reaction rates calculated from the fitting expressions and shown to be in a good agreement and the comparison of the two results are shown in Table (6).

Table 7 presents the comparison of thermonuclear reaction rates of some (α,n) medium elements reactions with other works as Roughton et.al. (Roughton *et al.*, 1983)

| Т9 (109 К) | CO | C1 | C2 |
|---------------|--------|--------|---------|
| 6 | 1.8494 | 0.9114 | -0.0281 |
| 6.25 | 3.9765 | 0.7593 | -0.0247 |
| 6.5 | 6.3605 | 0.5865 | -0.021 |
| 6.75 | 8.9003 | 0.4013 | -0.017 |
| 7 | 11.488 | 0.2126 | -0.013 |

Table 5. Free parameters C_i (C_0 , C_1 , and C_2) as a function of T_9 .

| 7.25 | 14.021 | 0.0288 | -0.0091 |
|------|--------|---------|---------|
| 7.5 | 16.418 | -0.1433 | -0.0056 |
| 7.75 | 18.627 | -0.2998 | -0.0023 |
| 8 | 20.644 | -0.4404 | 0.0006 |
| 8.25 | 22.523 | -0.5696 | 0.0032 |
| 8.5 | 24.388 | -0.6979 | 0.0058 |
| 8.75 | 26.451 | -0.8427 | 0.0086 |
| 9 | 29.02 | -1.0295 | 0.0124 |
| 9.25 | 32.512 | -1.2928 | 0.0176 |
| 9.5 | 37.472 | -1.6774 | 0.0253 |
| 9.75 | 44.58 | -2.2392 | 0.0366 |
| 10 | 54.667 | -3.0466 | 0.0528 |

Comparison between polynomial fitting expression (Best Fitting) of the adopted astrophysical S-Factor of (α, n) medium element reactions with those computed from Eq. (46).

| | 45 Sc(α ,n) | ^{48}V | ⁴⁸ Ti(a,r | n) ⁵¹ Cr | ⁵¹ V(α,n |) ⁵⁴ Mn | 55 Mn(α , | n) ⁵⁸ Co | ⁶² Ni(α,n |) ⁶⁵ Zn | ⁶⁶ Zn(α,n |) ⁶⁹ Ge |
|---------------|--|--|---|---|---|---|---|---|---|--|---|---|
| T9 (109 K) | ln[Na<σv> (cm3 s-1 mol- 1)] (Best Fitting) 4.04% | ln[Na< σv> (cm3 s- 1 mol- 1)] (Formu la) | ln[Na<σv> (cm3 s-1 mol-1)] (Best Fitting) 4.215% | ln[Na<σv > (cm3 s- 1 mol-1)] (Formula) | ln[Na<σv> (cm3 s-1 mol-1)] (Best Fitting) 3.801% | ln[Na<σv > (cm3 s- 1 mol-1)] (Formula) | ln[Na<σv> (cm3 s-1 mol-1)] (Best Fitting) 3.028% | ln[Na<σv > (cm3 s- 1 mol-1)] (Formula) | ln[Na<σv> (cm3 s-1 mol-1)] (Best Fitting) 3.785% | ln[Na<σ v> (cm3 s-1 mol- 1)] (Formul a) | ln[Na<σv> (cm3 s-1 mol- 1)] (Best Fitting) 9.071% | ln[Na<σv > (cm3 s- 1 mol-1)] (Formula) |
| 6 | 8.614±0.348 | 8.613 | 8.281±0.349 | 8.294 | 7.901±0.300 | 7.927 | 7.005±0.212 | 7.048 | 5.415±0.205 | 5.364 | 3.824±0.347 | 3.998 |
| 6.25 | 9.060±0.366 | 9.027 | 8.696±0.367 | 8.710 | 8.325±0.316 | 8.348 | 7.479±0.226 | 7.488 | 5.941±0.225 | 5.858 | 4.459±0.404 | 4.545 |
| 6.5 | 9.485±0.383 | 9.430 | 9.091±0.383 | 9.111 | 8.724±0.332 | 8.751 | 7.916±0.240 | 7.907 | 6.427±0.243 | 6.330 | 5.057±0.459 | 5.072 |
| 6.75 | 9.889±0.400 | 9.821 | 9.471±0.399 | 9.497 | 9.101±0.346 | 9.136 | 8.323±0.252 | 8.303 | 6.879±0.260 | 6.779 | 5.620±0.510 | 5.580 |
| 7 | 10.275±0.415 | 10.201 | 9.841±0.415 | 9.867 | 9.461±0.360 | 9.502 | 8.701±0.263 | 8.677 | 7.301±0.276 | 7.205 | 6.148±0.558 | 6.067 |
| 7.25 | 10.643±0.430 | 10.568 | 10.202±0.43 0 | 10.222 | 9.804±0.373 | 9.850 | 9.054±0.274 | 9.029 | 7.699±0.291 | 7.608 | 6.642±0.602 | 6.534 |
| 7.5 | 10.994±0.444 | 10.924 | 10.556±0.44 5 | 10.561 | 10.134±0.38 5 | 10.179 | 9.385±0.284 | 9.359 | 8.076±0.306 | 7.989 | 7.100±0.644 | 6.981 |
| 7.75 | 11.331±0.458 | 11.268 | 10.904±0.46 0 | 10.885 | 10.450±0.39 7 | 10.490 | 9.697±0.294 | 9.667 | 8.435±0.319 | 8.346 | 7.523±0.682 | 7.409 |
| 8 | 11.653±0.471 | 11.600 | 11.243±0.47 4 | 11.193 | 10.751±0.40 9 | 10.783 | 9.990±0.302 | 9.952 | 8.777±0.332 | 8.681 | 7.912±0.718 | 7.816 |
| 8.25 | 11.964±0.483 | 11.921 | 11.571±0.48 8 | 11.487 | 11.036±0.41 9 | 11.058 | 10.266±0.31 1 | 10.216 | 9.102±0.345 | 8.992 | 8.269±0.750 | 8.203 |
| 8.5 | 12.266±0.496 | 12.230 | 11.882±0.50 1 | 11.764 | 11.302±0.43 0 | 11.314 | 10.526±0.31 9 | 10.457 | 9.411±0.356 | 9.281 | 8.600±0.780 | 8.570 |
| 8.75 | 12.563±0.508 | 12.527 | 12.171±0.51 3 | 12.027 | 11.547±0.43 9 | 11.552 | 10.770±0.32 6 | 10.676 | 9.702±0.367 | 9.547 | 8.911±0.808 | 8.918 |
| 9 | 12.860±0.520 | 12.812 | 12.431±0.52 4 | 12.274 | 11.765±0.44 7 | 11.772 | 10.997±0.33 3 | 10.873 | 9.973±0.377 | 9.790 | 9.215±0.836 | 9.245 |
| 9.25 | 13.164±0.532 | 13.085 | 12.652±0.53 3 | 12.506 | 11.953±0.45 4 | 11.973 | 11.207±0.33 9 | 11.047 | 10.220±0.38 7 | 10.010 | 9.525±0.864 | 9.553 |
| 9.5 | 13.483±0.545 | 13.347 | 12.824±0.54 1 | 12.722 | 12.102±0.46 0 | 12.156 | 11.399±0.34 5 | 11.200 | 10.440±0.39 5 | 10.207 | 9.860±0.894 | 9.840 |
| 9.75 | 13.826±0.559 | 13.597 | 12.935±0.54 5 | 12.923 | 12.207±0.46 4 | 12.321 | 11.571±0.35 0 | 11.330 | 10.626±0.40 2 | 10.382 | 10.245±0.929 | 10.108 |
| 10 | 14.207±0.574 | 13.835 | 12.972±0.54 7 | 13.109 | 12.259±0.46 6 | 12.467 | 11.720±0.35 5 | 11.439 | 10.774±0.40 8 | 10.533 | 10.710±0.972 | 10.355 |

Table 7. comparison of the thermonuclear reaction rates in natural logarithm $\ln[Na < \sigma v > (cm^3 s^{-1} mol^{-1})]$ of some (α ,n) medium element reactions with other works.

| | ⁴⁸ Ti(α,n | 51 Cr | $^{51}V(\alpha,n)^{54}Mn$ | | ${}^{50}\mathrm{Cr}(\alpha,n){}^{53}\mathrm{Fe}$ | | 55 Mn(α ,n) 58 Co | |
|---------------|----------------------------|-----------------|----------------------------|-----------------|--|-----------------|-------------------------------------|-----------------|
| T9 (109 K) | Roughton et al. 1983 | Present Work | Roughton et al. 1983 | Present Work | Roughton et al. 1983 | Present Work | Roughton et al. 1983 | Present Work |
| 2 | -10.054 | -11.904 | -10.680 | -9.599 | -15.936 | -16.080 | -12.652 | -13.626 |

| 3 | -1.347 | -2.179 | -2.120 | -1.796 | -5.655 | -5.707 | -3.507 | -3.953 |
|---|---|--|---|--|---|--|--|--|
| 4 | 3.401 | 2.918 | 2.708 | 2.723 | -0.211 | -0.181 | 1.629 | 1.283 |
| 5 | 6.446 | 6.098 | 5.914 | 5.780 | 3.258 | 3.313 | 5.011 | 4.654 |
| 6 | 8.556 | 8.290 | 8.189 | 7.973 | 5.670 | 5.728 | 7.378 | 7.015 |
| 7 | 10.127 | 9.894 | 9.903 | 9.607 | 7.496 | 7.490 | 9.116 | 8.755 |
| 8 | 11.327 | 11.117 | 11.184 | 10.860 | 8.882 | 8.826 | 10.491 | 10.081 |
| 9 | 12.301 | 12.076 | 12.206 | 11.845 | 9.999 | 9.869 | 11.608 | 11.120 |
| 10 | 13.102 | 12.846 | 13.060 | 12.636 | 10.897 | 10.700 | 12.506 | 11.952 |
| | ⁵⁴ Fe(a,n |) ⁵⁷ Ni | ⁵⁹ Co(a | n) ⁶² Cu | ⁶³ Cu(a | n) ⁶⁶ Ga | ⁶⁶ Zn(a | n) ⁶⁹ Ge |
| T9 | Roughton et | D | Roughton | | Roughton | _ | Roughton et | |
| (109 K) | al. 1983 | Work Work | et al. 1983 | Present Work | et al. 1983 | Present Work | al. 1983 | Present Work |
| (109 K) | al. 1983 -18.526 | Vork | et al. 1983 -16.811 | Present Work -17.051 | et al. 1983 -25.945 | Present Work -25.653 | al. 1983 -25.759 | Present Work -26.138 |
| (109 K) 2 3 | al. 1983 -18.526 -6.908 | Present Work -18.474 -6.867 | et al. 1983 -16.811 -5.991 | Present Work -17.051 -6.195 | et al. 1983 -25.945 -10.871 | Present Work -25.653 -11.147 | al. 1983 -25.759 -10.820 | Present Work -26.138 -11.238 |
| (109 K) 2 3 4 | al. 1983 -18.526 -6.908 -0.968 | Present Work -18.474 -6.867 -0.942 | et al. 1983 -16.811 -5.991 -0.174 | Present Work -17.051 -6.195 -0.380 | et al. 1983 -25.945 -10.871 -3.270 | Present Work -25.653 -11.147 -3.901 | al. 1983 -25.759 -10.820 -3.244 | Present Work -26.138 -11.238 -3.708 |
| (109 K) 2 3 4 5 | al. 1983 -18.526 -6.908 -0.968 2.708 | Present Work -18.474 -6.867 -0.942 2.699 | et al. 1983 -16.811 -5.991 -0.174 3.526 | Present Work -17.051 -6.195 -0.380 3.311 | et al. 1983 -25.945 -10.871 -3.270 1.335 | Present Work -25.653 -11.147 -3.901 0.427 | al. 1983 -25.759 -10.820 -3.244 1.386 | Present Work -26.138 -11.238 -3.708 0.826 |
| (109 K) 2 3 4 5 6 | al. 1983 -18.526 -6.908 -0.968 2.708 5.193 | Present Work -18.474 -6.867 -0.942 2.699 5.179 | et al. 1983 -16.811 -5.991 -0.174 3.526 6.131 | Present Work -17.051 -6.195 -0.380 3.311 5.865 | et al. 1983 -25.945 -10.871 -3.270 1.335 4.431 | Present Work -25.653 -11.147 -3.901 0.427 3.291 | al. 1983 -25.759 -10.820 -3.244 1.386 4.522 | Present Work -26.138 -11.238 -3.708 0.826 3.841 |
| (109 K) 2 3 4 5 6 7 | al. 1983 -18.526 -6.908 -0.968 2.708 5.193 7.003 | Present Work -18.474 -6.867 -0.942 2.699 5.179 6.978 | et al. 1983 -16.811 -5.991 -0.174 3.526 6.131 8.039 | Present Work -17.051 -6.195 -0.380 3.311 5.865 7.726 | et al. 1983 -25.945 -10.871 -3.270 1.335 4.431 6.659 | Present Work -25.653 -11.147 -3.901 0.427 3.291 5.317 | al. 1983 -25.759 -10.820 -3.244 1.386 4.522 6.791 | Present Work -26.138 -11.238 -3.708 0.826 3.841 5.979 |
| (109 K) 2 3 4 5 6 7 8 | al. 1983 -18.526 -6.908 -0.968 2.708 5.193 7.003 8.434 | Present Work -18.474 -6.867 -0.942 2.699 5.179 6.978 8.339 | et al. 1983 -16.811 -5.991 -0.174 3.526 6.131 8.039 9.547 | Present Work -17.051 -6.195 -0.380 3.311 5.865 7.726 9.134 | et al. 1983 -25.945 -10.871 -3.270 1.335 4.431 6.659 8.343 | Present Work -25.653 -11.147 -3.901 0.427 3.291 5.317 6.821 | al. 1983 -25.759 -10.820 -3.244 1.386 4.522 6.791 8.517 | Present Work -26.138 -11.238 -3.708 0.826 3.841 5.979 7.568 |
| (109 K) 2 3 4 5 6 7 8 9 | al. 1983 -18.526 -6.908 -0.968 2.708 5.193 7.003 8.434 9.547 | Present Work -18.474 -6.867 -0.942 2.699 5.179 6.978 8.339 9.401 | et al. 1983 -16.811 -5.991 -0.174 3.526 6.131 8.039 9.547 10.714 | Present Work -17.051 -6.195 -0.380 3.311 5.865 7.726 9.134 10.231 | et al. 1983 -25.945 -10.871 -3.270 1.335 4.431 6.659 8.343 9.680 | Present Work -25.653 -11.147 -3.901 0.427 3.291 5.317 6.821 7.975 | al. 1983 -25.759 -10.820 -3.244 1.386 4.522 6.791 8.517 9.852 | Present Work -26.138 -11.238 -3.708 0.826 3.841 5.979 7.568 8.789 |



Fig. 3. The variation of the natural logarithm of the thermonuclear reaction rates with the atomic number Z for the $^{45}Sc(\alpha,n)^{48}V$, $^{48}Ti(\alpha,n)^{51}Cr$, $^{51}V(\alpha,n)^{54}Mn$, $^{55}Mn(\alpha,n)^{58}Co$, $^{62}Ni(\alpha,n)^{65}Zn$, and $^{66}Zn(\alpha,n)^{69}Ge$ reactions at fixed values of T₉.









(c)

Fig. 4. C_i coefficients against T₉, for C_0 , C_1 , and C_2 respectively. The solid line represents the fitted curve through the data.

5. Conclusions

- 1-The astrophysical S-factor, S(E), was starting with an increase and then decreased irregularly by increasing the center of mass energy, this because of Coulomb barrier penetration $exp(2\pi\eta)$.
- 2-The astrophysical S-factor increased with increasing atomic number Z of target nuclei at a fixed center of mass energy.
- 3-The thermonuclear reaction rates, $N_A < \sigma v >$, were increased with increasing T₉ because by increasing the T₉ the charged interacting particles need to overcome the existing Coulomb barrier.
- 4-The thermonuclear reaction rates decreased with increasing atomic number Z of target nuclei at fixed T₉ because as Z increased Coulomb barrier increased.
- 5-The astrophysical S-factor and Thermonuclear reaction rates calculated in the present work are in good agreement with those measured previously by other works.

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